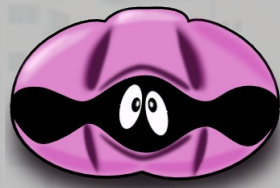


Many-Body Perturbation theory: Basic concepts and approximations

Andrea Marini

October 21, Barranquilla, Colombia

FLASH **it**



www.yambo-code.eu



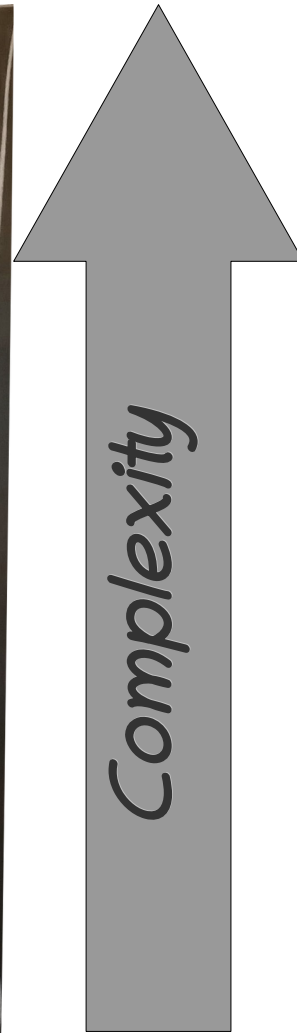
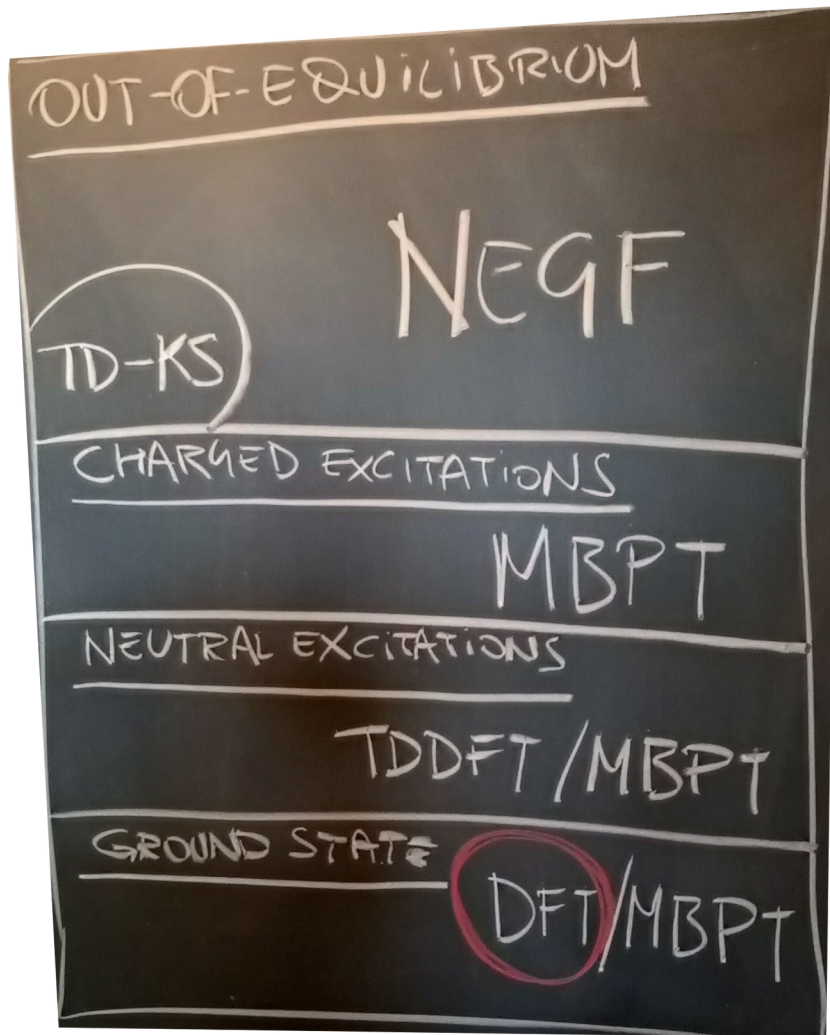
Istituto di Struttura
della Materia



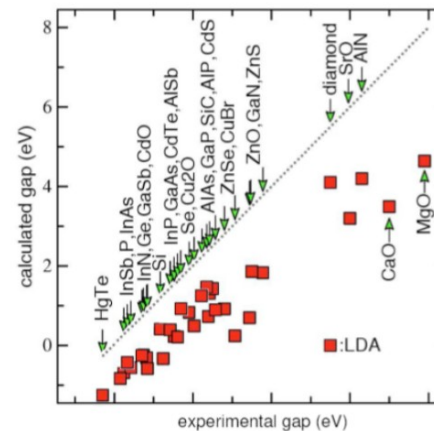
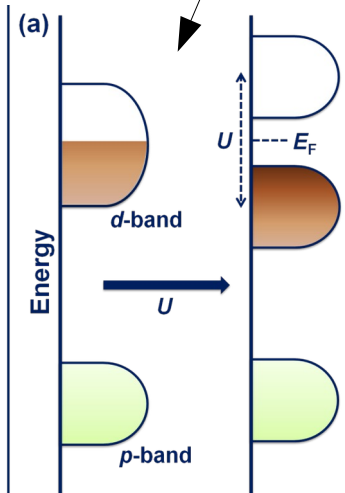
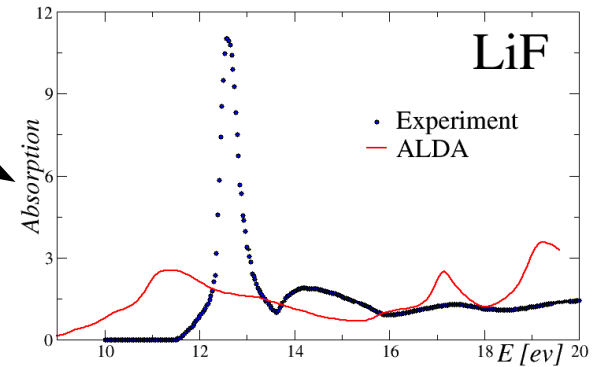
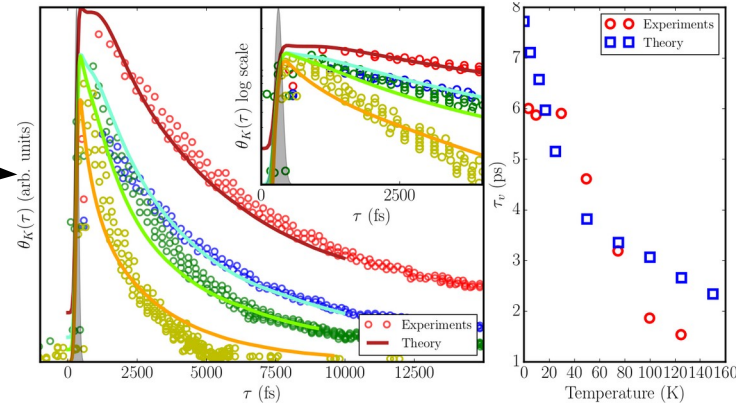
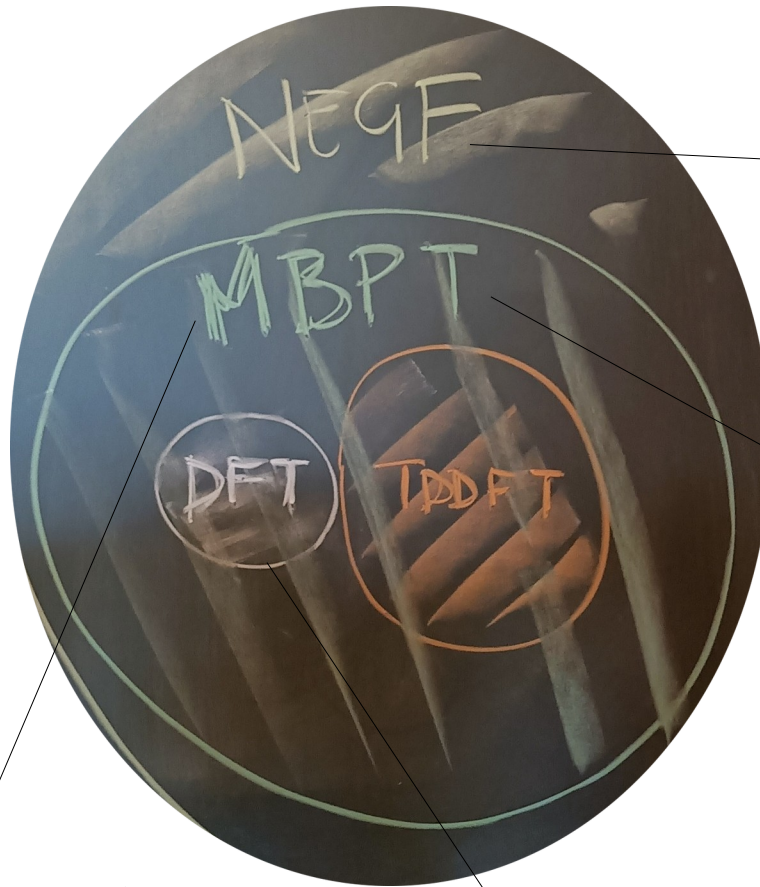
Ultrafast Science Laboratory of the
Material Science Institute National Research Council
(Monterotondo Stazione, Italy)

<http://www.yambo-code.eu/andrea>

Different physics, different approaches



Different physics, different approaches



Si:
0.47 eV (LDA) vs 1.1 eV (expt)

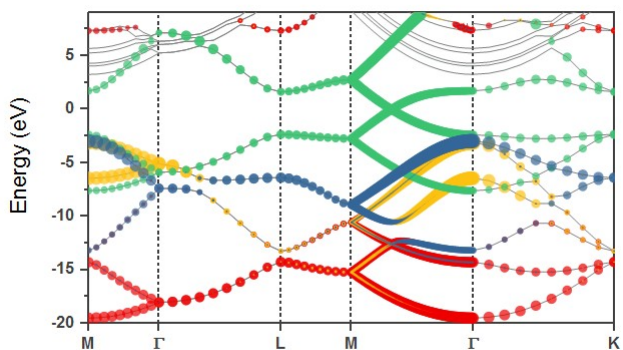
GaAs:
0.30 eV (LDA) vs 1.4 eV (expt)

The Many-Body problem

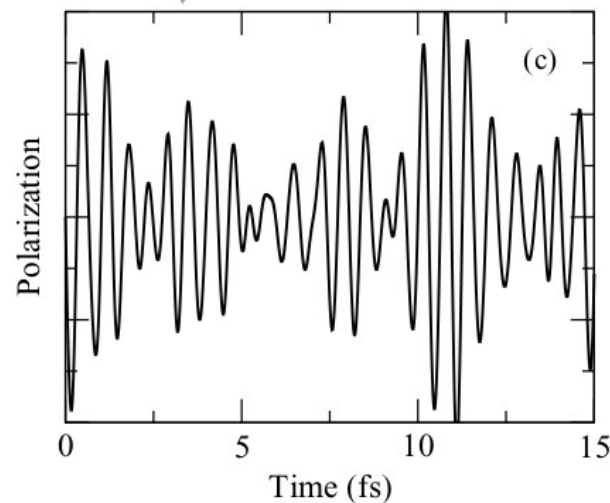
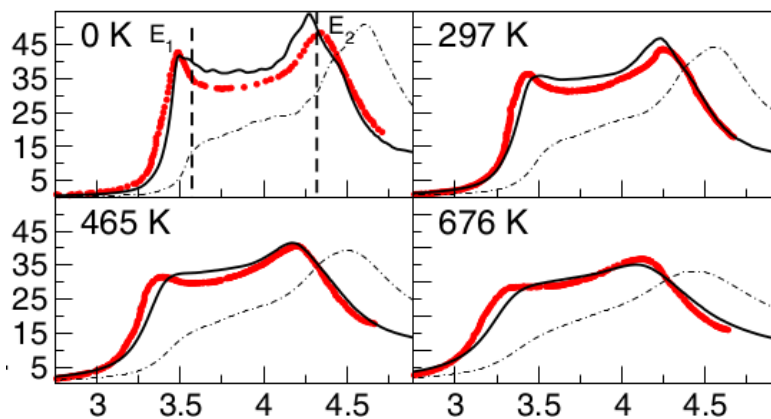
$$H = \sum_i h(x_i, p_i) + \frac{1}{2} \sum_{i \neq j} |x_i - x_j|^{-1}$$



The Many-Body Problem: a micro-macro connection



$$\hat{H} = \sum_i \hat{h}(x_i) + \frac{1}{2} \sum_{ij}' \frac{1}{|x_i - x_j|}$$

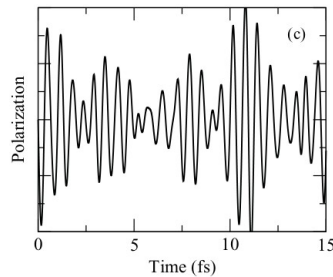


A (very) hard job!

$$\langle N | = \overline{(|N\rangle)}$$

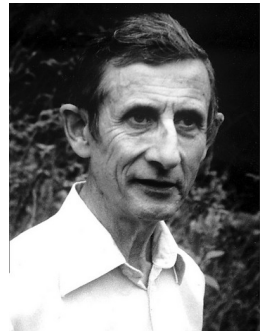
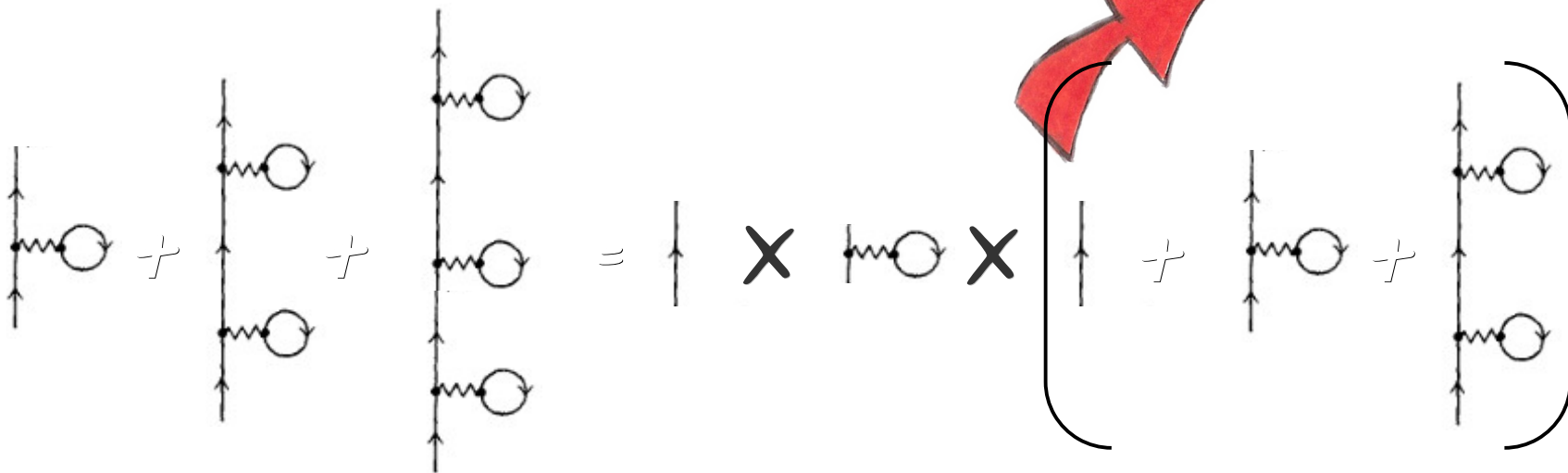
QM

$$A = \langle N | \hat{A} | N \rangle$$



$$|N(t)\rangle = U(t, t_0) |N(t_0)\rangle$$

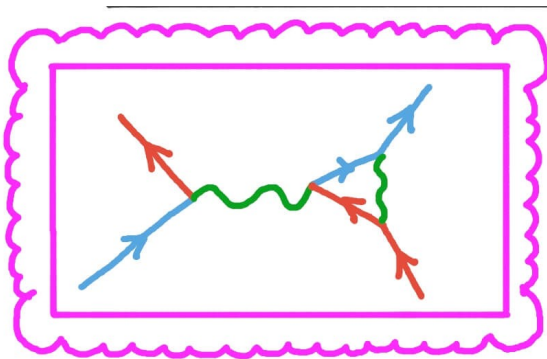
Diagrams



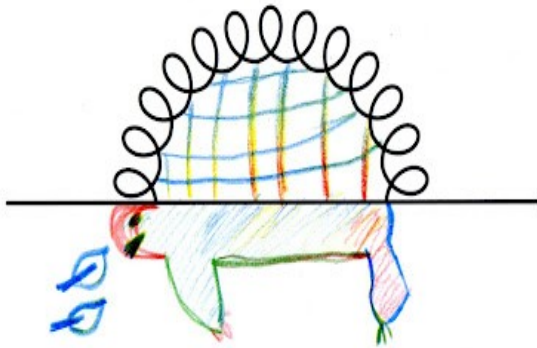
Outline



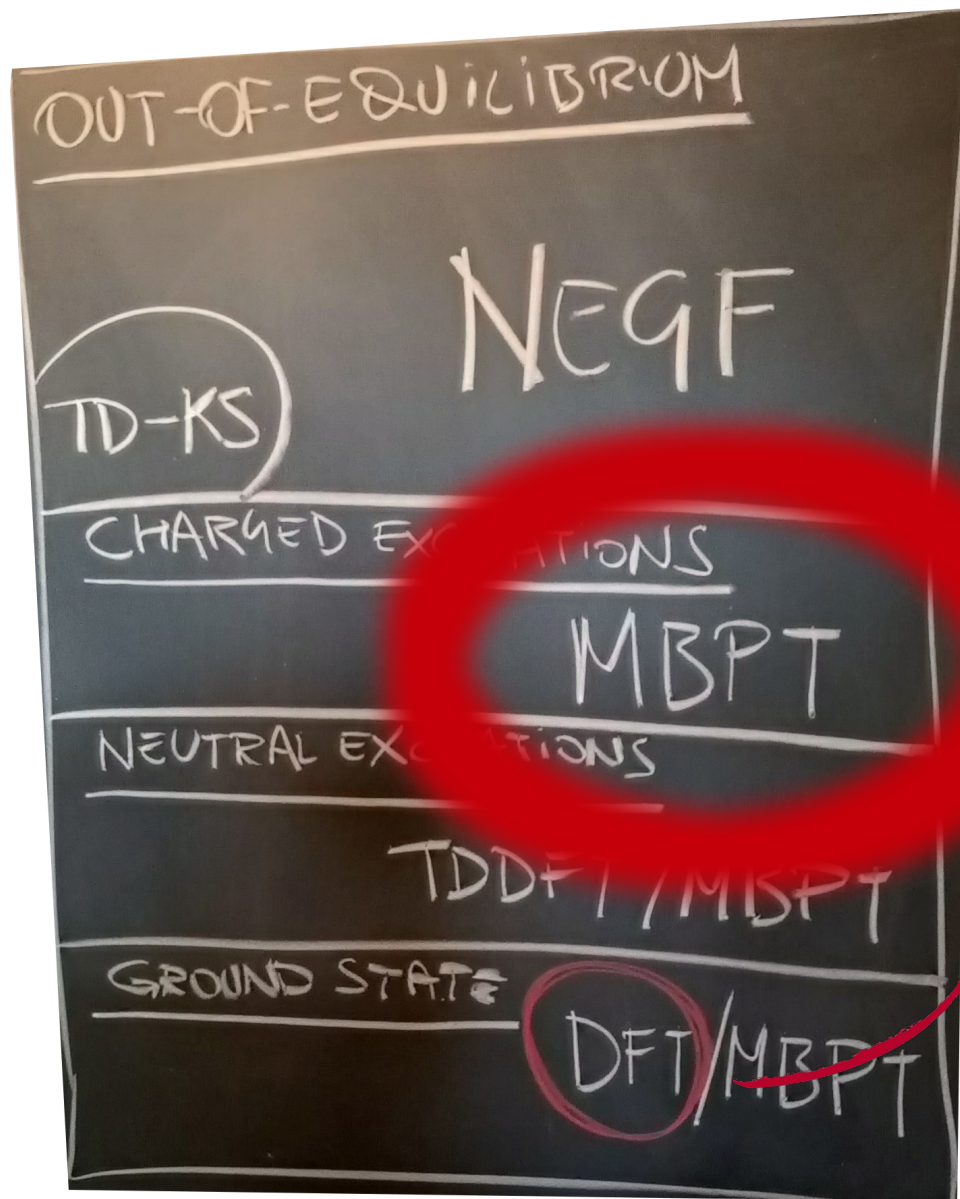
Many-Body Perturbation Theory for dummies



Feynman diagrams for dummies



The “zoo” of diagrammatic approximations

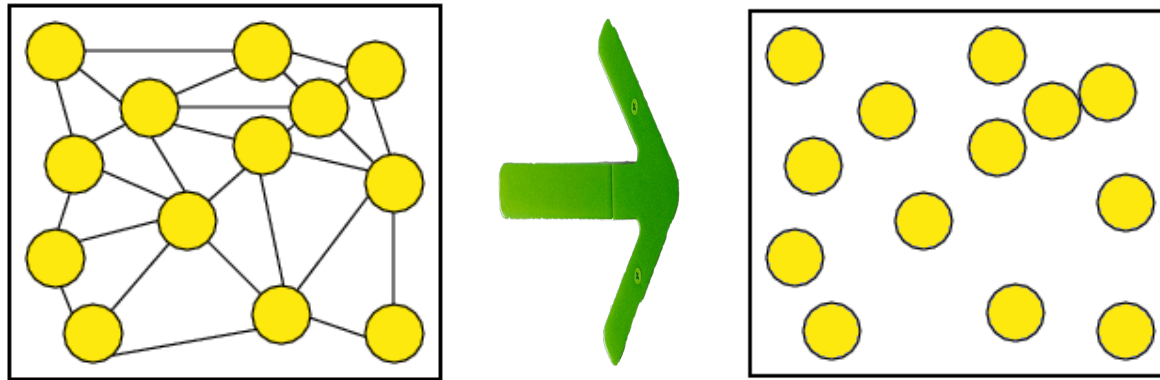
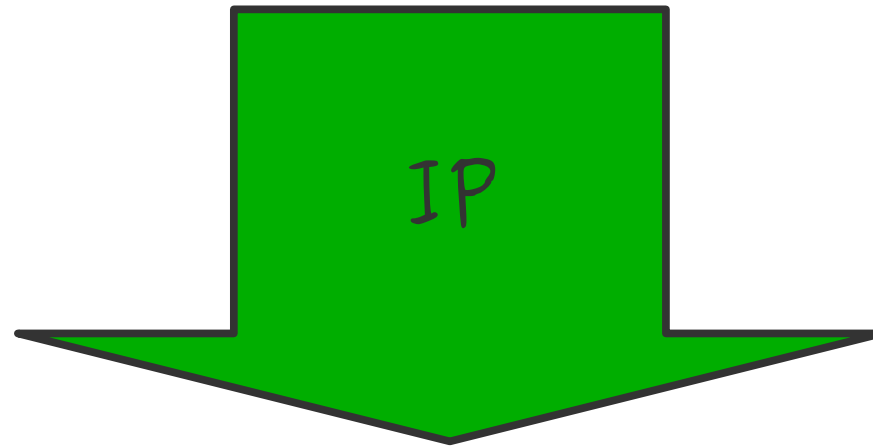


Many-Body Perturbation
Theory for dummies



The Many-Body problem

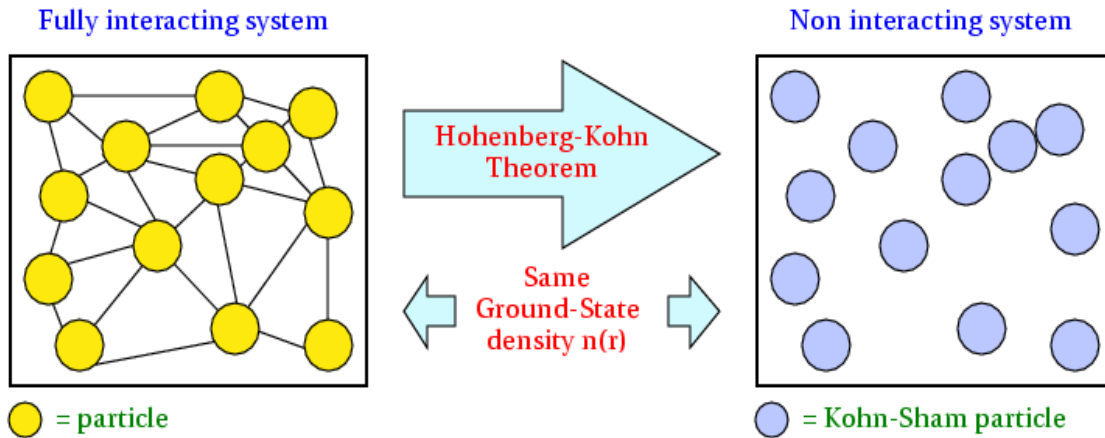
$$H = \sum_i h(x_i, p_i) + \frac{1}{2} \sum_{i \neq j} |x_i - x_j|^{-1}$$



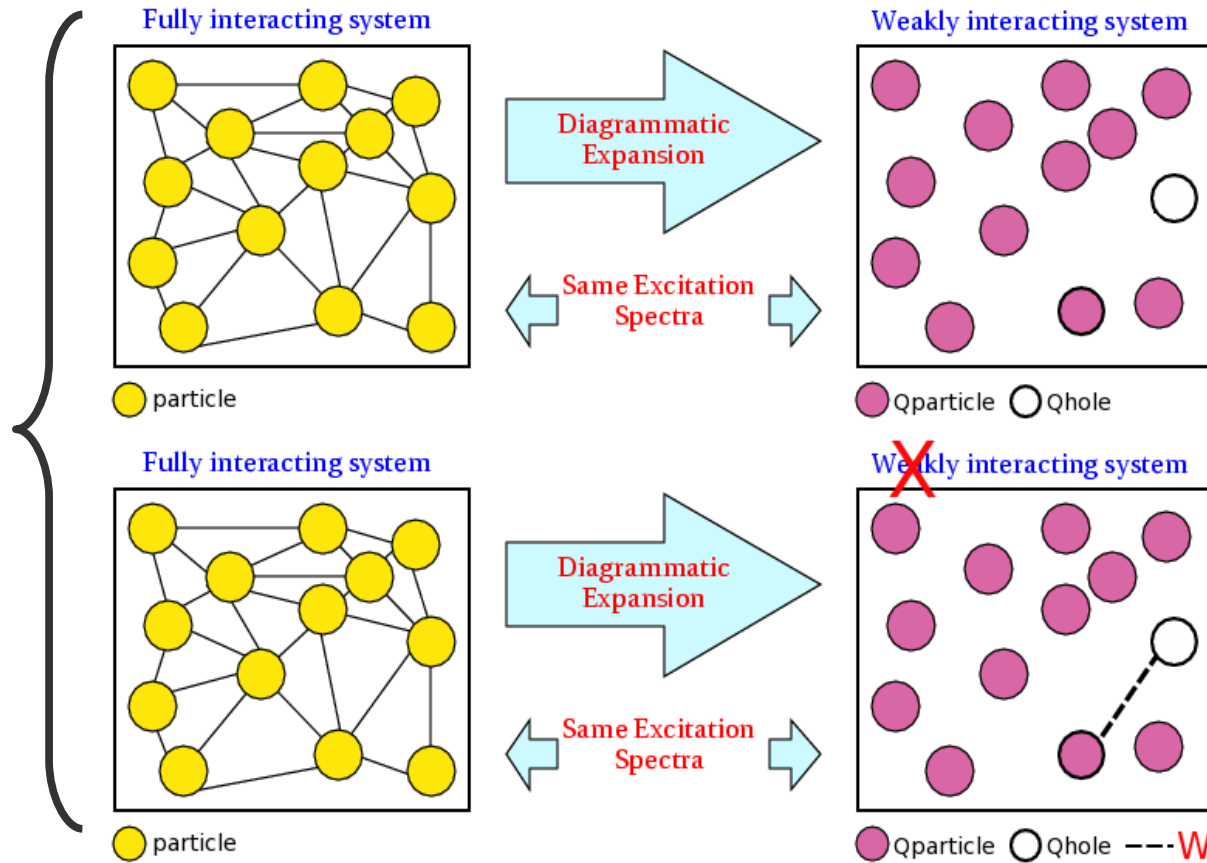
$$H \approx \sum_i h(x_i)$$

The Many-Body problem

DFT



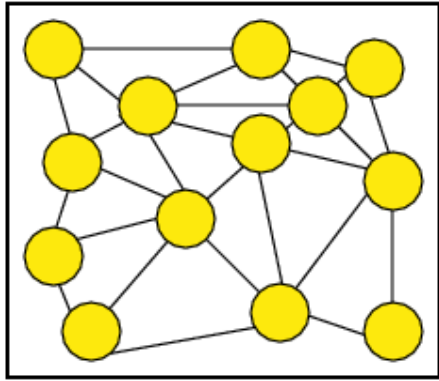
MBPT



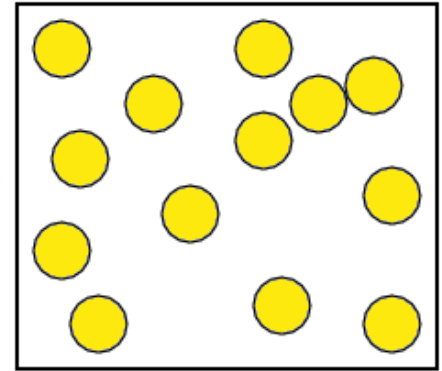
The Many-Body problem: 1 particle approx

$$H = \sum_i h(x_i) + \frac{1}{2} \sum_{i \neq j} \frac{1}{|x_i - x_j|^{-1}}$$

$$H = \sum_i h(x_i)$$



$$\hat{h}|n\rangle = \epsilon_n|n\rangle$$



$$|N_0\rangle = \prod_{n \in \text{filled}} |n\rangle$$

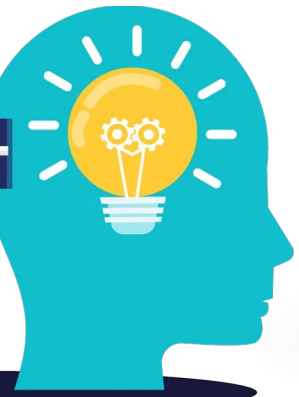
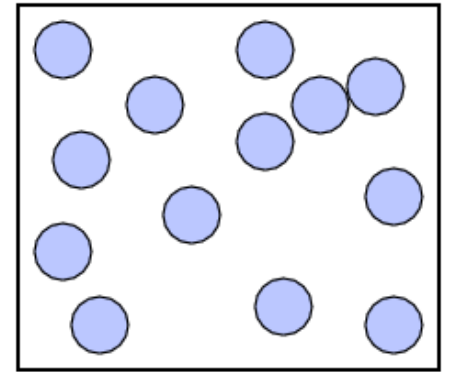
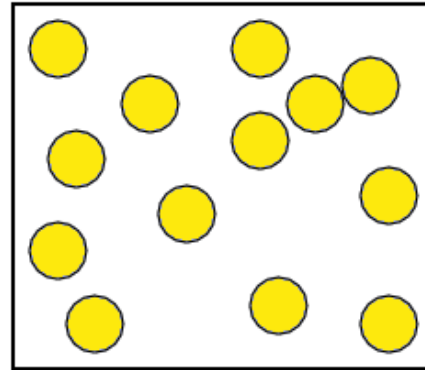
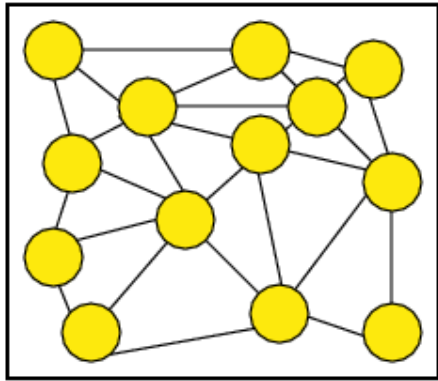


$$\langle N | \hat{A} | N \rangle \approx F_N [\{A_n\}]$$

$$\langle N_0 | \hat{H} | N_0 \rangle = \sum_{n \in \text{filled}} \epsilon_n$$

Quasiparticles...

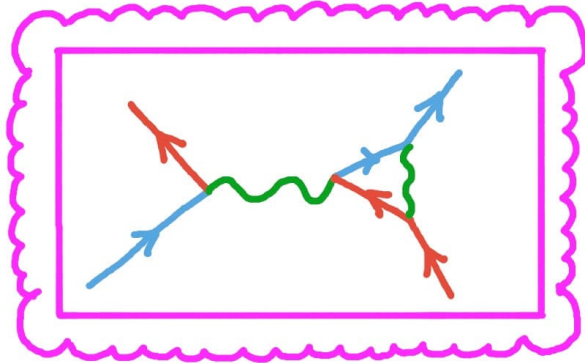
$$H = \overbrace{\sum_i h(x_i)}^{\hat{H}_0} + \frac{1}{2} \sum_{i \neq j} |x_i - x_j|^{-1} = h + H'$$



For MBPT KS is
a mean-field
quasiparticle

The goal of the Many Body methods
is to rewrite the fully interacting
problem as an as much
independent as possible counter-
part

FEYNMAN DIAGRAMS



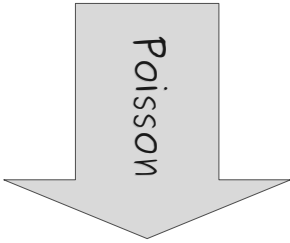
... a beautifully rendered,
pictorial representation
of the great physicist
Richard P. Feynman...

Feynmann diagrams for
dummies

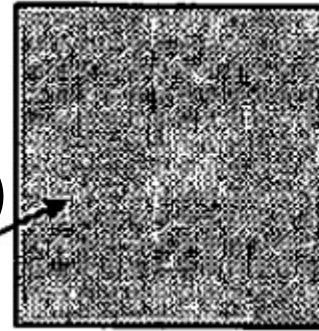


The Coulomb interaction

$$\nabla^2 V(\mathbf{x}, t) = 4\pi n(\mathbf{x}, t)$$



$$n(\mathbf{x}, t)$$

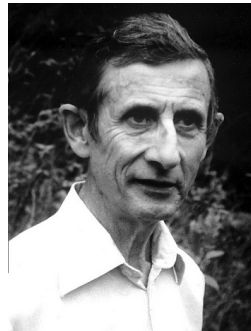


Let's add an external (oscillating) charge in the system

$$V(\mathbf{r}, t) = \int d\mathbf{r}' \frac{n(\mathbf{r}', t)}{|\mathbf{r} - \mathbf{r}'|}$$

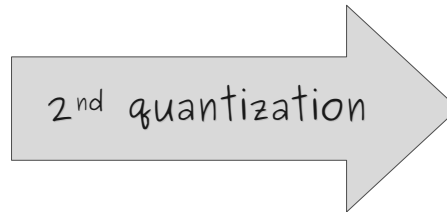


$$\hat{\psi}(\mathbf{r}) = \sum_{\mathbf{k}} \phi_{\mathbf{k}}(\mathbf{r}) \hat{d}_{\mathbf{k}}$$



$$\hat{n}(\mathbf{r}) = \hat{\psi}^\dagger(\mathbf{r}) \hat{\psi}(\mathbf{r})$$

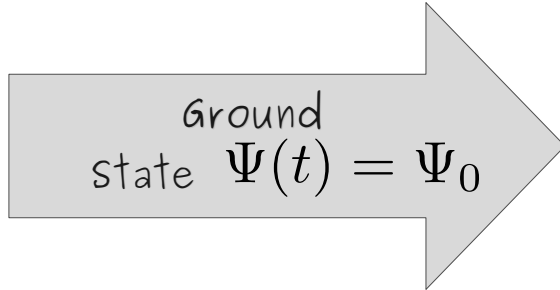
2nd quantization



$$\hat{H}(t) = \hat{H}_0 + \int \hat{n}(\mathbf{r}) V(\mathbf{r}, t)$$

The time-dependent, interacting density (Kubo)

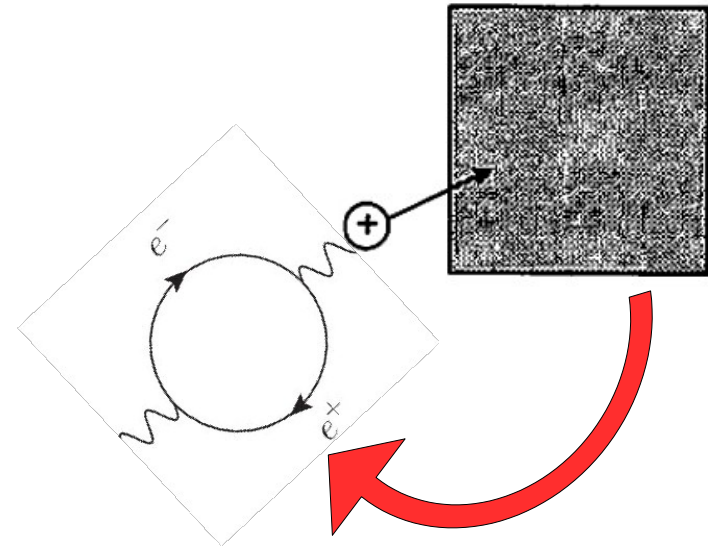
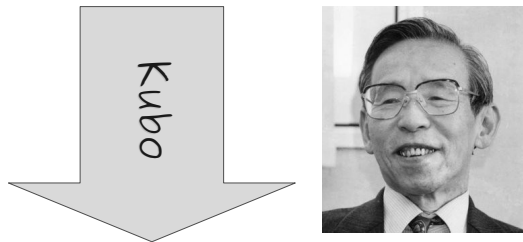
$$n(\mathbf{r}, t) = \langle \Psi(t) | \hat{\psi}^\dagger(\mathbf{r}) \hat{\psi}(\mathbf{r}) | \Psi(t) \rangle$$



Density Functional Theory

Semi-classical excitation

$$\hat{H}(t) = \hat{H}_0 + \int \hat{n}(\mathbf{r}) V(\mathbf{r}, t)$$

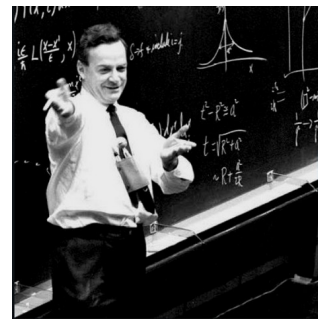


$$n(\mathbf{r}, t) = n_0(\mathbf{r}) + \int d\mathbf{r}' \int_0^t dt' \chi^R(\mathbf{r}t, \mathbf{r}'t') V(\mathbf{r}', t')$$

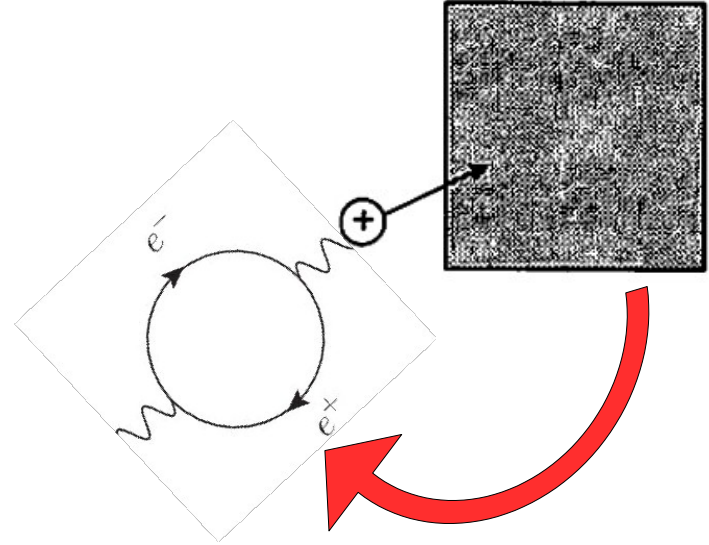
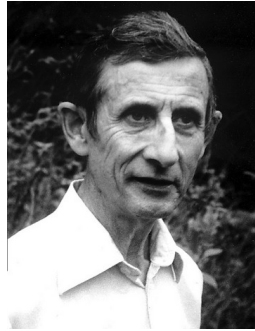
$$\chi(\mathbf{r}t, \mathbf{r}'t') = -i \langle \Psi | [\hat{n}(\mathbf{r}, t), \hat{n}(\mathbf{r}', t')] | \Psi \rangle \theta(t - t')$$

The response (Green's) function

$$n(\mathbf{r}, t) = n_0(\mathbf{r}) + \int d\mathbf{r}' \int_0^t dt' \underbrace{\chi^R(\mathbf{r}t, \mathbf{r}'t')} V(\mathbf{r}', t')$$

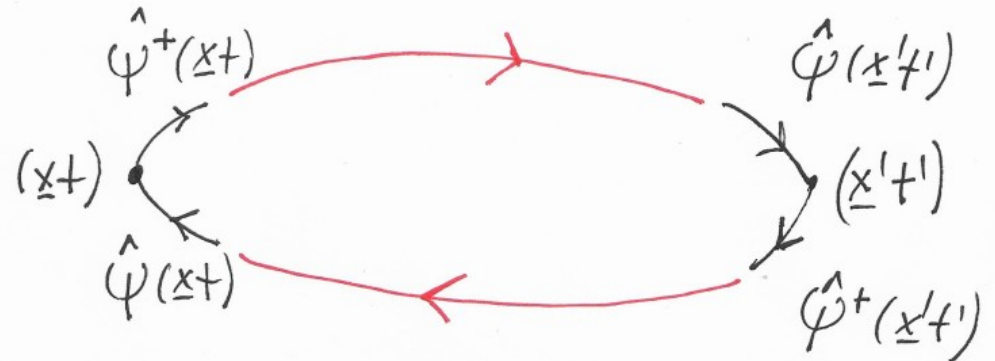


$$\hat{\psi}(\mathbf{r}) = \sum_{\mathbf{k}} \phi_{\mathbf{k}}(\mathbf{r}) \hat{d}_{\mathbf{k}}$$



$$\hat{\psi}^\dagger(\mathbf{r}t) \longrightarrow \hat{\psi}(\mathbf{r}'t')$$

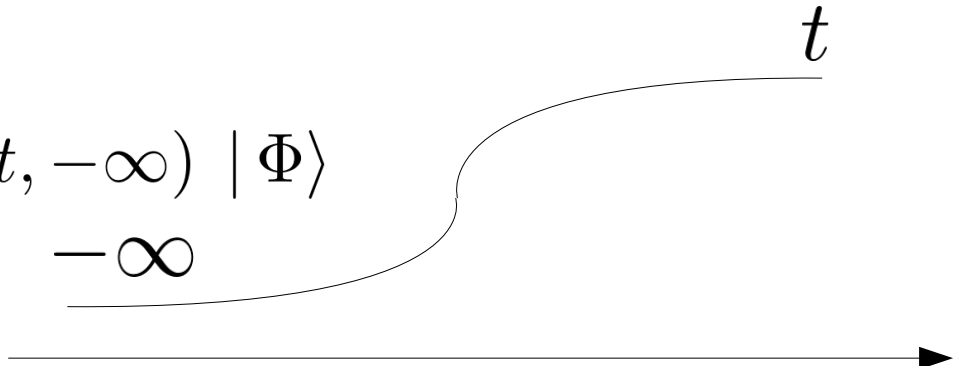
$$\langle \Psi | [\hat{n}(\mathbf{r}, t), \hat{n}(\mathbf{r}', t')] | \Psi \rangle$$

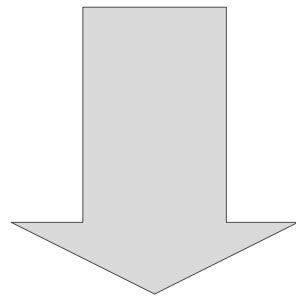


The time-dependent, interacting density

$$n(\mathbf{r}, t) = \left\langle \Psi(t) \left| \hat{\psi}^\dagger(\mathbf{r}) \hat{\psi}(\mathbf{r}) \right| \Psi(t) \right\rangle$$

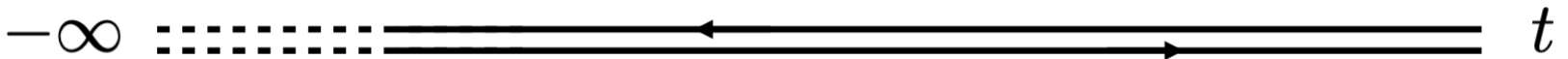
$$|\Psi(t)\rangle = \hat{U}(t, t_0) |\Psi(t_0)\rangle \longrightarrow \hat{U}(t, -\infty) |\Phi\rangle$$

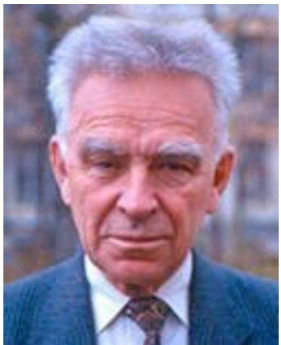




$$\hat{U}(t) \equiv \hat{U}(t, -\infty)$$

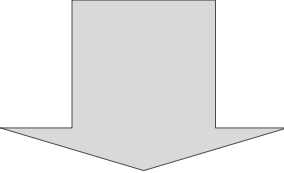
$$n(\mathbf{r}, t) = \sum_{\mathbf{k}\mathbf{k}'} \phi_{\mathbf{k}}^*(\mathbf{r}) \phi_{\mathbf{k}'}(\mathbf{r}) \left\langle \Psi(t) \left| \hat{d}_{\mathbf{k}}^\dagger \hat{U}(t) \hat{U}^\dagger(t) \hat{d}_{\mathbf{k}'} \right| \Psi(t) \right\rangle$$



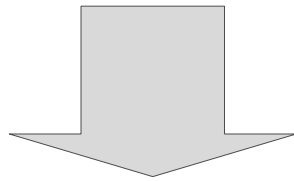


The evolution operator (scattering potential)

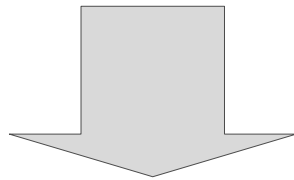
$$\hat{H}(t) = \hat{H}_0 + \underbrace{\hat{V}(t)}$$


$$\int \hat{n}(\mathbf{r}) V(\mathbf{r}, t)$$

$$i \frac{d}{dt} \hat{U}_0(t) = \hat{H}_0 \hat{U}_0(t)$$



$$\hat{U}(t) = \hat{U}_0(t) \hat{F}(t)$$



$$\hat{F}(t) = 1 - i \int_{-\infty}^t dt_1 \hat{V}_I(t_1) + \underbrace{(-i)^2 \int_{-\infty}^t dt_1 \int_{-\infty}^{t_1} dt_2 \hat{V}_I(t_1) \hat{V}_I(t_2)} + \dots$$

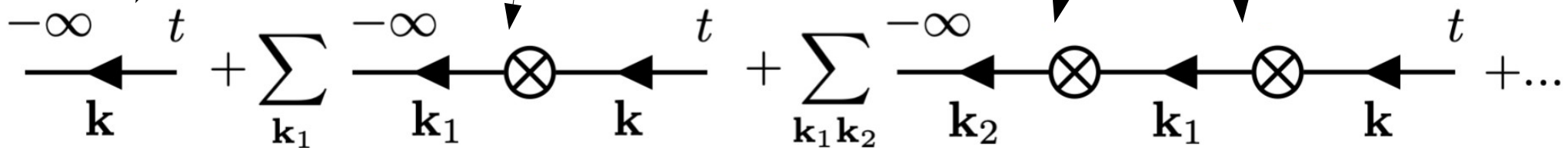
Constrained time integrals
 $t > t_1 > t_2 > \dots$

Half the dynamics...

$$\langle \Psi(t) | \hat{d}_{\mathbf{k}}^\dagger \hat{U}(t) \hat{U}^\dagger(t) \hat{d}_{\mathbf{k}'} | \Psi(t) \rangle = \delta_{\mathbf{k}\mathbf{k}'} - \langle \Psi(t) | \hat{d}_{\mathbf{k}} \hat{U}(t) \hat{U}^\dagger(t) \hat{d}_{\mathbf{k}'}^\dagger | \Psi(t) \rangle$$

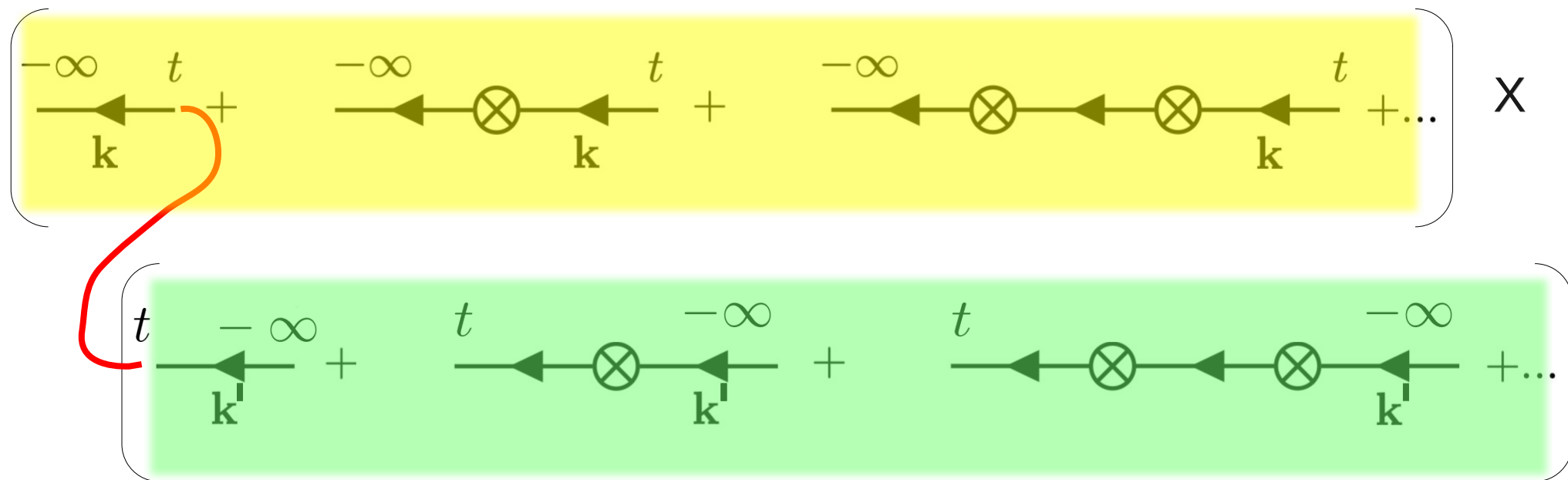
$$\hat{U}^\dagger(t) \hat{d}_{\mathbf{k}}^\dagger | \Psi(t) \rangle = \underbrace{F^\dagger(t)}_{\hat{F}^\dagger(t)} \hat{U}_0^\dagger(t) \hat{d}_{\mathbf{k}}^\dagger | \Psi(t) \rangle$$

$$\hat{F}^\dagger(t) = 1 + i \int_{-\infty}^t dt_1 \hat{V}_I(t_1) + i^2 \int_{-\infty}^t dt_1 \int_{-\infty}^{t_1} dt_2 \hat{V}_I(t_2) \hat{V}_I(t_1) + \dots$$



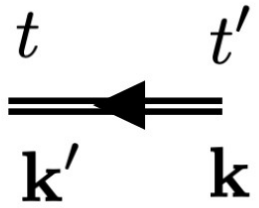
Green's Functions

$$\left\langle \Psi(t) \left| \hat{d}_{\mathbf{k}}^\dagger \hat{U}(t) \hat{U}^\dagger(t) \hat{d}_{\mathbf{k}'}^\dagger \right| \Psi(t) \right\rangle =$$



$$= \begin{array}{c} t & t \\ \leftarrow & \leftarrow \\ \mathbf{k}' & \mathbf{k} \end{array} + \begin{array}{c} t & t_1 & t \\ \leftarrow & \otimes & \leftarrow \\ \mathbf{k}' & & \mathbf{k} \end{array} + \begin{array}{c} t & t_2 & t_1 & t \\ \leftarrow & \otimes & \leftarrow & \otimes & \leftarrow \\ \mathbf{k}' & & & & \mathbf{k} \end{array} + \dots$$

The Dyson equation



$$= \begin{array}{c} t \quad t' \\ \leftarrow \quad \rightarrow \\ \hline k' \quad k \end{array} + \begin{array}{c} t \quad t_2 \quad t_1 \quad t' \\ \leftarrow \quad \bullet \quad \leftarrow \quad \rightarrow \\ \hline k' \quad \Sigma \quad k \end{array} + \begin{array}{c} t \quad t_2 \quad t_1 \quad t_3 \quad t_4 \quad t' \\ \leftarrow \quad \bullet \quad \leftarrow \quad \bullet \quad \leftarrow \quad \rightarrow \\ \hline k' \quad \Sigma \quad G_0 \quad \Sigma \quad k \end{array} + \dots$$

$$= \begin{array}{c} t \quad t' \\ \leftarrow \quad \rightarrow \\ \hline k' \quad k \end{array} + \begin{array}{c} t \quad t_2 \quad t_1 \quad t' \\ \leftarrow \quad \bullet \quad \leftarrow \quad \rightarrow \\ \hline k' \quad \Sigma \quad G \end{array}$$

The Dyson Equation



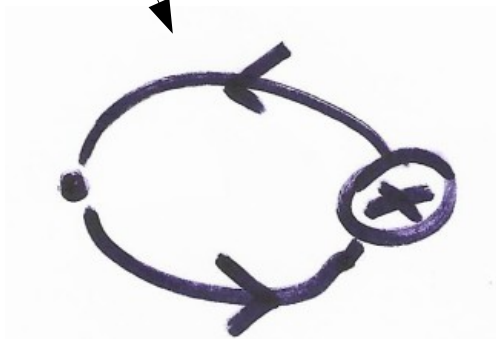
Green's Functions: Kubo revisited

$$\langle \Psi(t) | \hat{d}_{\mathbf{k}}^\dagger \hat{U}(t) \hat{U}^\dagger(t) \hat{d}_{\mathbf{k}'}^\dagger | \Psi(t) \rangle =$$

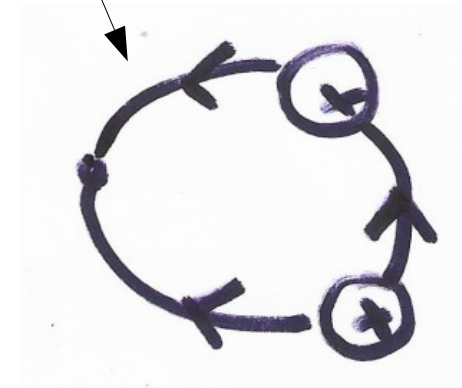
$$= \begin{array}{c} t \quad t \\ \leftarrow \quad \leftarrow \\ \mathbf{k}' \quad \mathbf{k} \end{array} + \begin{array}{c} t \quad t_1 \quad t \\ \leftarrow \quad \otimes \quad \leftarrow \\ \mathbf{k}' \quad \quad \mathbf{k} \end{array} + \underbrace{\begin{array}{c} t \quad t_2 \quad t_1 \quad t \\ \leftarrow \quad \otimes \quad \leftarrow \quad \otimes \quad \leftarrow \\ \mathbf{k}' \quad \quad \quad \mathbf{k} \end{array}} + \dots$$



$n_0(\mathbf{r})$



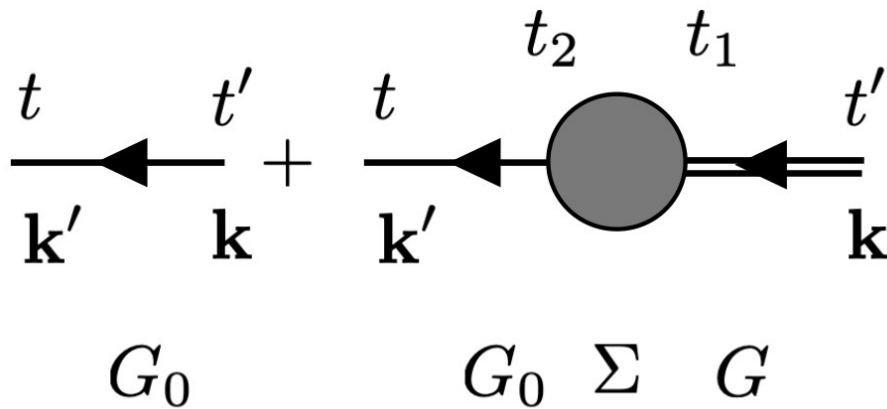
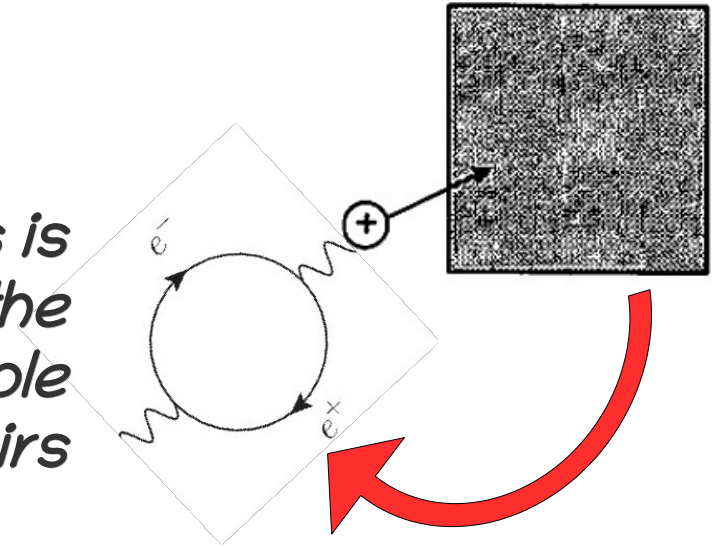
$$\int d\mathbf{r}' \int_0^t dt' \chi^R(\mathbf{r}t, \mathbf{r}'t') V(\mathbf{r}', t')$$



Non-linear terms

Key messages

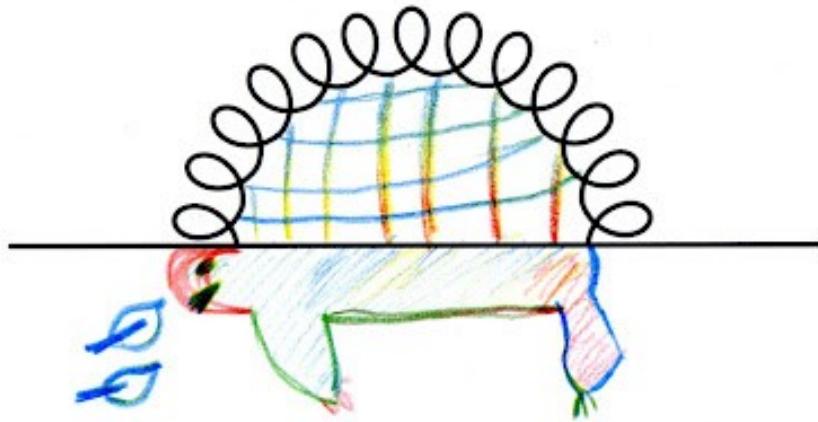
Basic MBPT process is screening through the excitation of electron-hole (neutral) pairs



The very same process can be easily described by using a diagrammatic representation

MBPT is (by far) more powerful when we move to more complicated interaction potentials





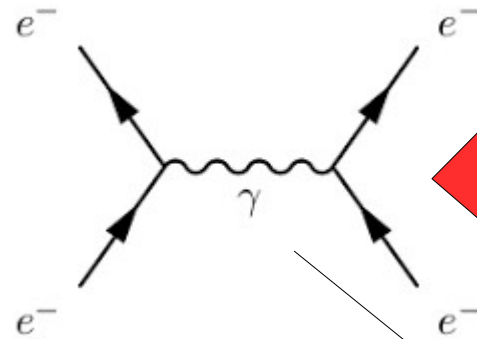
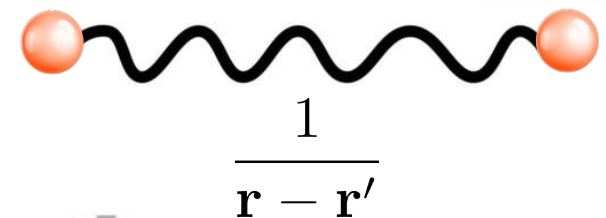
The “zoo” of diagrammatic approximations



The Coulomb interaction (revisited)

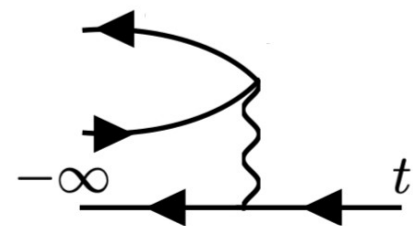
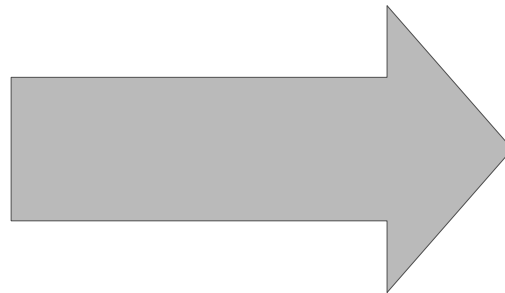
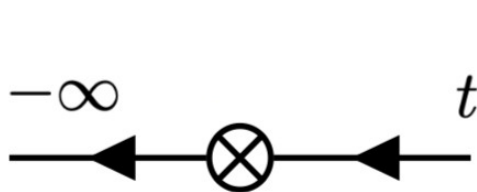


$$V(\mathbf{r}, t) = \int d\mathbf{r}' \frac{n(\mathbf{r}', t)}{|\mathbf{r} - \mathbf{r}'|}$$



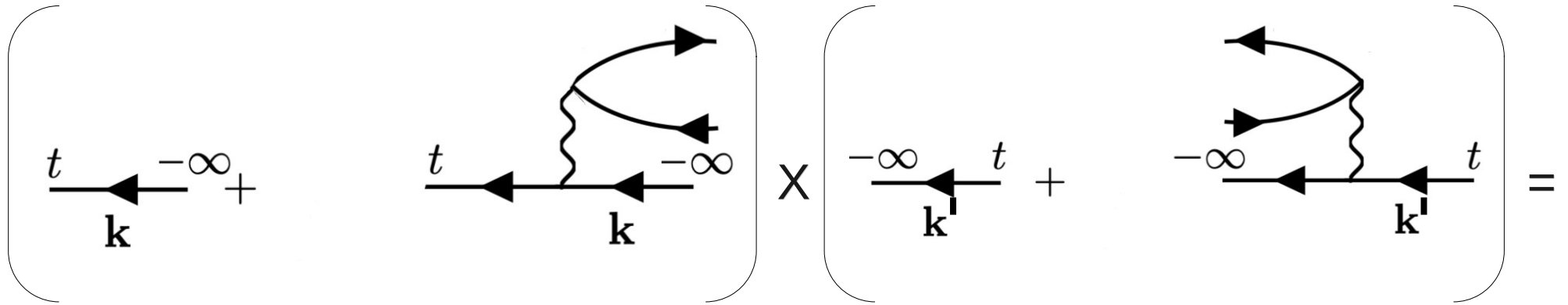
$$\hat{U}(t) = \hat{U}_0(t) \hat{F}(t)$$

$$\hat{F}(t) = 1 - i \int_{-\infty}^t dt_1 \hat{V}_I(t_1) + (-i)^2 \int_{-\infty}^t dt_1 \int_{-\infty}^{t_1} dt_2 \hat{V}_I(t_1) \hat{V}_I(t_2) + \dots$$

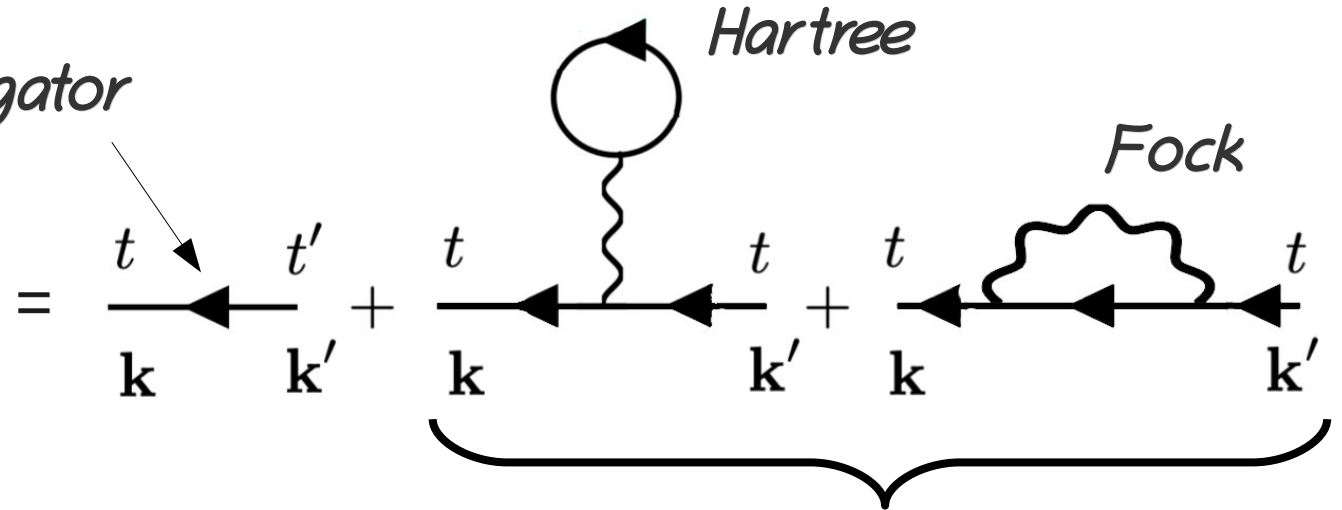


Feynman diagrams in the fully interacting case

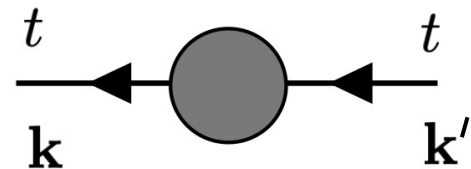
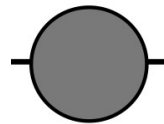
$$\langle \Psi(t) | \hat{d}_{\mathbf{k}}^\dagger \hat{U}(t) \hat{U}^\dagger(t) \hat{d}_{\mathbf{k}'}^\dagger | \Psi(t) \rangle =$$



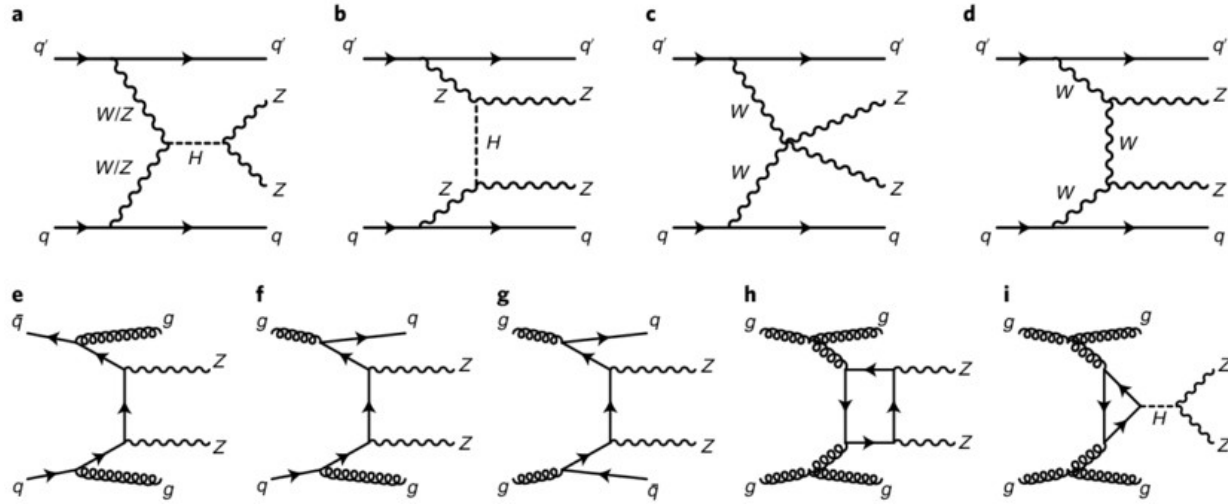
The Propagator



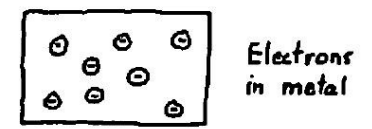
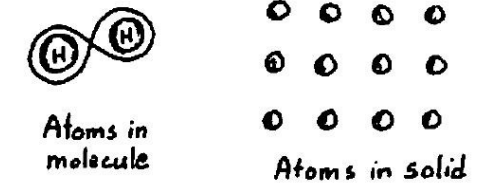
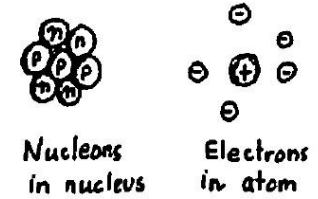
The Self-Energy



Feynman diagrams in the fully interacting case



Use Physical arguments to choose specific classes of diagrams !!!



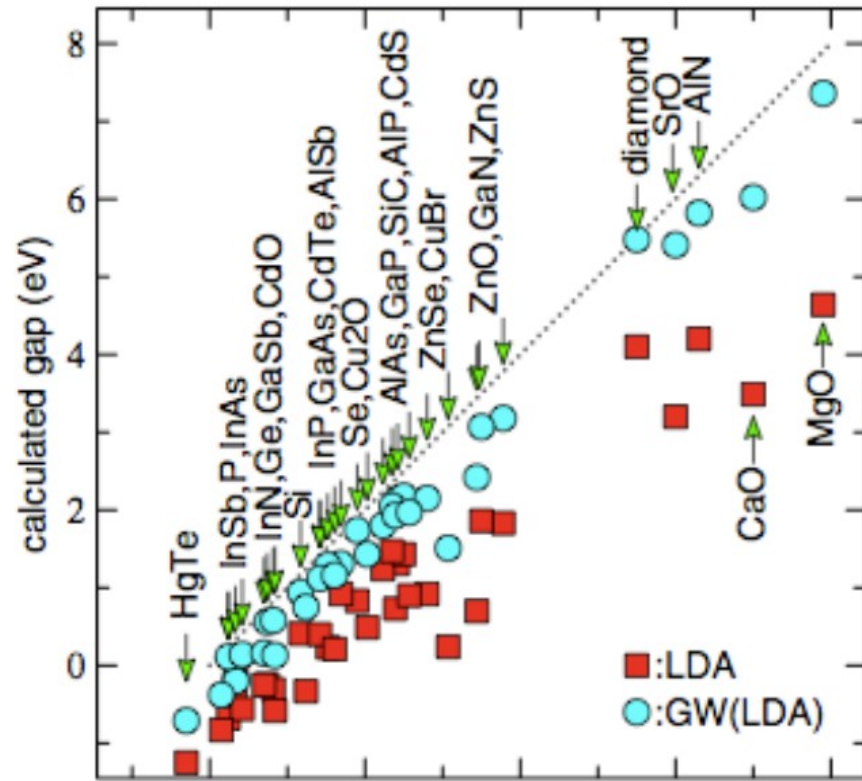
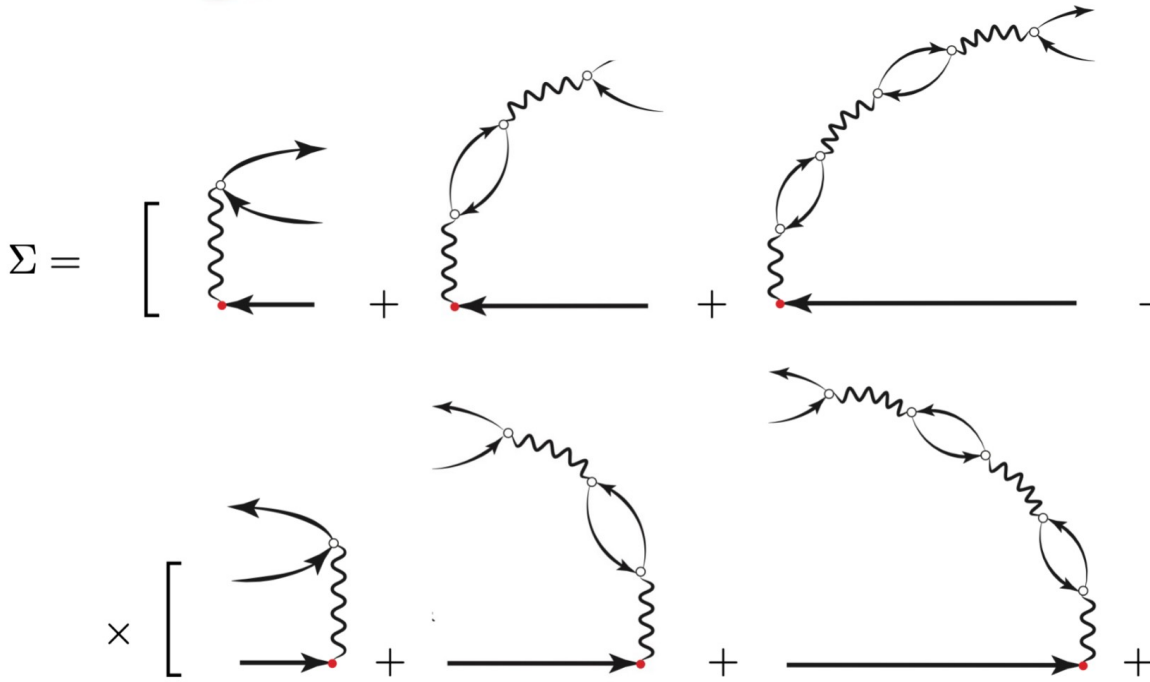
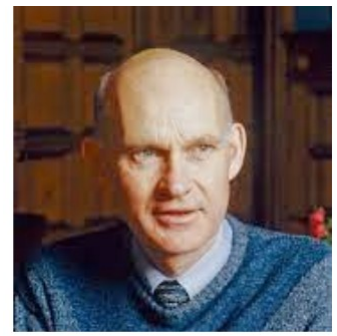
Short-range interactions ?

High density regime ?

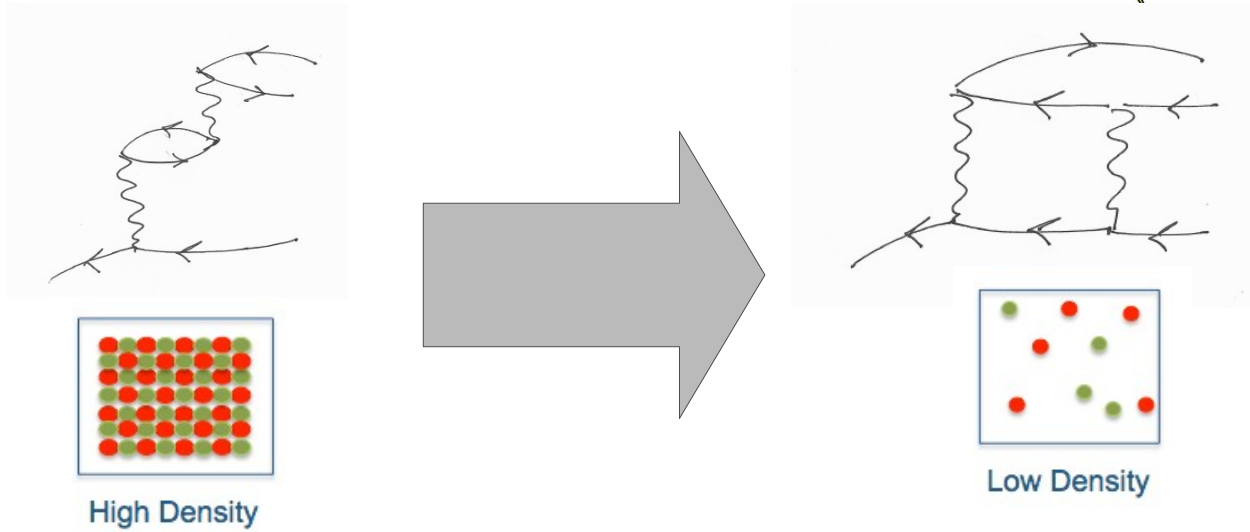
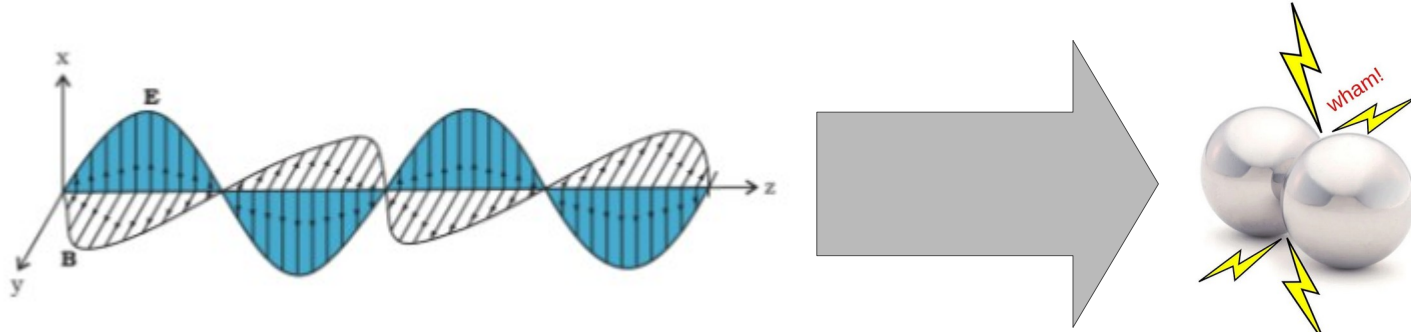
Low density regime ?

Conserving approximations

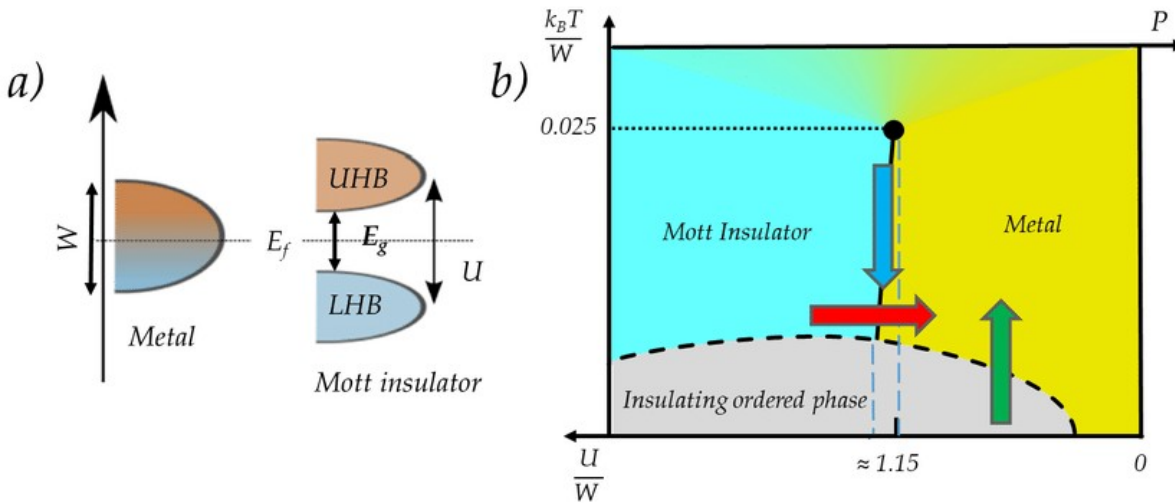
The GW approximation



The T-matrix approximation



VIKTOR MIKHAĬLOVICH
GALITSKIĬ
(1924-1981)

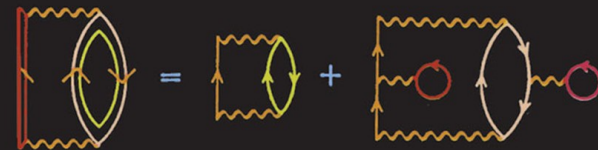


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Richard D. Mattuck



A Guide to
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- [Energy Loss Spectroscopy](#) , F. Sottile

Many-body Theory

- [PhD lectures: MBPT and Yambo](#) , L. Chiodo et al.
- [Introduction to Many Body Physics](#) , Piers Coleman
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